

## The BEC-BCS crossover and highly excited exciton-polariton condensates

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6. Negative mas dispersion due to quantum depletion

Assume particles at condensate scatter into the positive and negative branches to conserve energy and momentum.

Applying the parametric theory [1]

$$i\hbar \frac{d}{dt} \begin{pmatrix} \hat{p}_{k}(t) \\ \hat{p}_{-k}^{\dagger}(t)e^{-i2\varepsilon_{0}t/\hbar} \end{pmatrix} = \hat{M}_{k}^{par} \begin{pmatrix} \hat{p}_{k}(t) \\ \hat{p}_{-k}^{\dagger}(t)e^{-i2\varepsilon_{0}t/\hbar} \end{pmatrix} + \begin{pmatrix} \hat{F}_{k}(t) \\ -\hat{F}_{-k}^{\dagger}(t)e^{-i2\varepsilon_{0}t/\hbar} \end{pmatrix}$$



Wavenumber

Condensate energy

Where  

$$\hat{M}_{k}^{par} = \begin{pmatrix} E_{k}(n) + U(n) - i\gamma & U(n) \\ -U(n) & 2\varepsilon_{0} - E_{k}(n) - U(n) - i\gamma \end{pmatrix}$$

Photoluminescence is then defined by  $PL(k,t,\omega) \propto \left|C_k\right|^2 \operatorname{Re} \int_{0}^{+\infty} d\tau e^{-i(\omega-i0^+)\tau} \left\langle \hat{p}_k^{\dagger}(t+\tau) \hat{p}_k(t) \right\rangle$ Here, we use parameters from BCS model [2] in parametric

[1] C. Ciuti, P. Schwendimann, A. Quattropani, PRB 63, 041303(R) (2001)

[2] T. Byrnes, T. Horikiri. N. Ishida, Y. Yamamoto, PRL 105, 186402 (2010)

