Non-equilibrium many-body physics with alkaline earth atoms and polar molecules

Ana Maria Rey





\$ Funding \$

NSF, AFOSR, ARO, ARO-DARPA-OLE,

The 11th US-Japan Joint Seminar 2013

"Ultimate Quantum Systems of Light and Matter- Control and Applications"

The Sr team:



Jun Ye

M. Swallows, M. Martin, M. Bishof, S. Blatt, X. Zhang, C. Benko

Theory team:



A. Gorshkov, M. Foss-Feig, K. Hazzard, B. Zhu, S. Manmana, M. Lukin

KRb team:









B. Yan, S. Moses, J. Covey, B. Neyenhuis and B. Gadway



Ultra-cold Matter

Fully controllable quantum systems

The most precise measurements,

e.g. clocks

Quantum sensors

A tool for understanding quantum complexity







Many body physics with clocks

Understanding quantum systems from few- to many-body with "clock" precision and control



Alkaline earth - super coherence





JILA Ultra-coherent spectroscopy: Nicholson *et al.,* Phys. Rev. Lett. 109 (2012) 230801

Ye's talk this afternoon Q~10¹⁵, seconds coherence time



Magic wave length

Ye, Kimble, & Katori, Science **320**, 1734 (2008).



No Doppler, No Recoil No Stark shift



What happens in the real experiment with N particles?

S = N/2

Non-interacting: Collective-spin

$$S_{x,y,z} = \sum_{n} S_n^{x,y,z}$$
 Interactions?

³P₀(e) ¹D lattice clock: ^(a) T~µK

 ${}^{1}S_{0}(g)$

000

Effective spin 1/2 system during clock interrogation

Dominant p-wave collisions

Both elastic and inelastic

Mode occupation is conserved. No laser/interaction induced mode changing collisions.

Decoupled motional/spin





$$S_{x,y,z} = \sum_{n} S_{n}^{x,y,z}$$



Same Hamiltonian that two component Bose Einstein Condensate: Sorensen, Moller, Cirac, Zoller, Lewenstein, ... >



A. M. Rey, L. Jiang, M. Fleischhauer, E. Demler, and M. D. Lukin, PRA 77, 052305 (2008).



Treat other surrounding atoms as an average





$H = - \Box S^z)^2 \rightarrow 2 \Box S^z \Box B S^z$



$B=-2 \square S^{z} \square - N \bigcirc \theta$

Spin precesses with a modified rate with depends on atom number



 θ controlled by first pulse

M. Martin et al arXiv:1212.6291

Quantum correlations – beyond mean field

Quantum correlations should manifest on the amplitude of the oscillations

- At the mean field level interactions only affect the precession rate.
- Amplitude remains constant

But..... in the experiments there are many pancakes with different atom number. Due to interactions the pancakes with more atoms precess faster.



Signal adds → amplitude of the oscillations decay due to dephasing or destructive interference between pancakes

• Atom number decay aslo leads to decay of the amplitude

M. Martin et al arXiv:1212.6291

Comparisons with experiment

Ramsey fringe decay vs. the spin tipping angle

To eliminate the effect of decay we normalize the amplitude with atom number $2\pi/3$



M. Martin et al arXiv:1212.6291

Quantum fluctuations – beyond mean field

Interplay between interactions and decoherence: complicated



Quantum correlations induce faster decay of the amplitude

We were able to solve the full master Eq for the collective model.



Polar molecules





Gorshkov et al: PRL.107.115301(2011), PRA 84,033619 (2011)





dipole moment Select two dressed levels : Effective spin ¹/₂ system

Related previous work

Other schemes: *Micheli* et al, *Nat. Phys.* **2** *341* (2006); *Brennen* et al, *NJP* **9** *138* (2007); *Buechler* et al, *Nat. Phys.* **3** *726* (2007); *Perez-Rios*, et al *NJP* **12**, *103007*; *Wall-Carr Phys. Rev. A 82*, *013611* (2010)...

Spin model: Frozen molecules in a lattice $H_{dd} = d_i d_j V_{dd}^{ij} \qquad V_{dd}^{ij} = \frac{(1 - 3\cos^2 \theta)}{|r_i - r_j|^3}$ d₁ $|\uparrow\rangle = |N = 1, M = 0\rangle$ • Project d_i on the two selected rotational levels $\left|\downarrow\right\rangle = \left|N=0, M=0\right\rangle$ $d_{i} = \hat{z} \sum_{\sigma \sigma'} d_{\sigma,\sigma'} |\sigma_{i}\rangle \langle \sigma_{i}'| d_{\sigma\sigma'} = \langle \sigma | d | \sigma' \rangle$ $H_{dd} = \sum V_{dd}^{ij} \left[J_z(d_{\sigma\sigma'}) S_i^z S_j^z + J_\perp(d_{\sigma\sigma'}) \left(S_i^x S_j^x + S_i^y S_j^y \right) \right]$ Ising Flip-flop $J_{\perp} = 2(d_{\uparrow \perp})^2$ 0.6 J_{\perp} L/Jch $J_z = (d_{\uparrow\uparrow} - d_{\downarrow\downarrow})^2$ J_z 0.10.00.20.40.6 $J_{ch} \sim 2\pi \times 250 \text{Hz} \rightarrow 1\text{ms}$ $\sim E$ d_0/d_p 4 kV/cm

Quantum Magnetism

□Use direct dipole-dipole interaction to generate direct strong (~KHz) spin exchange interaction:

10-100 larger than super-exchange or magnetic dipoles



Fully tunable coefficients by E field (microwaves) dipole moment Gorshkov et al: PRL.107.115301(2011), PRA 84,033619 (2011)

Long-range (1/r³) and anisotropic interactions:
S. R. Manmana *et al* PRB 87, 081106(R) (2013),
A. V. Gorshkov *et al* arXiv:1301.5636

Spin *temperature*, not motional *temperature* matters:
Relevant ratio is interaction time (~ms) to cloud lifetime (25 sec!):
K. R.A. Hazzard et al PRL 110, 075301 (2013)



- Empty sites act as defects
- Need to perform disorder average

Dipolar interactions will be visible in the Ramsey fringe contrast even in dilute samples



Full solution 1D-DMRG

K. R.A. Hazzard et al PRL 110, 075301 (2013)



3D-KRb lattice experiment

Current experiments are carried out in a 3D lattice with a B field

B: determines quantization axis

| **□ | =** |N=1,M=-1□

I ____ |N=0,M=0 ___

Polarization trapping light



Magic wavelength for their lattice

B. Neyenhuis *et al* Phys. Rev. Lett. 109, 230403 (2012)



• Non-trivial dependence on the geometry due to the anisotropic dipolar interactions.

$$V_{dd}^{ij} = \frac{(1 - 3\cos^2 \theta)}{|r_i - r_j|^3}$$

3D-lattice experiment in B field $H = J_{\perp} \sum_{\langle i,j \rangle} V_{ij} [S^{x}{}_{i}S^{x}{}_{j} + S^{y}{}_{i}S^{y}{}_{j}] \qquad J_{\perp} = -(d_{\uparrow\downarrow})^{2}$

Cluster Expansion



- Spins grouped in cluster of max size g.
- Intra-cluster interactions kept
- Inter-cluster interactions neglected or treated as a perturbation.

Preliminary comparisons with experiment



Dynamical decoupling

Learn from NMR: By applying the proper pulse sequence it is possible to eliminate dipolar interactions.

 $\left(\frac{\pi}{2}\right)_{x}^{(\pi)_{\chi}} \left(\frac{\pi}{2}\right)_{-\chi} \left(\frac{\pi}{2}\right)_{\chi} \left(\frac{\pi}{2}\right)_{\phi}$

Room-Temperature Quantum Bit Memory Exceeding One Second



Conclusions

- Ultra-cold polar matter offers a unique controllable laboratory for the exploration of many-body physics
- Strongly interacting open driven quantum systems
- Manifestation of quantum magnetism observable even in a nonquantum degenerate gas

• Rich physics a lot to be understood

Thanks