

$$\begin{aligned}
|C(a, b) - C(a, b')| &\leq \int |A(a, \lambda)[B(b, \lambda) - B(b', \lambda)]\rho(\lambda)d\lambda \\
&\quad \vdots \quad \leftarrow |A(a, \lambda)| = 1 \quad \quad \quad \nearrow \text{positive value} \\
&= \int |B(b, \lambda) - B(b', \lambda)|\rho(\lambda)d\lambda \quad \text{--- (1)}
\end{aligned}$$

$$\begin{aligned}
|C(a', b) + C(a', b')| &\leq \int |A(a', \lambda)[B(b, \lambda) + B(b', \lambda)]\rho(\lambda)d\lambda \\
&= \int |B(b, \lambda) + B(b', \lambda)|\rho(\lambda)d\lambda \quad \text{--- (2)}
\end{aligned}$$

(1) + (2) \implies

$$\begin{aligned}
&|C(a, b) - C(a, b')| + |C(a', b) + C(a', b')| \\
&\leq \int \underbrace{[|B(b, \lambda) - B(b', \lambda)| + |B(b, \lambda) + B(b', \lambda)|]}_{= 2 \text{ for all combinations of } [A(a), B(b)] = (1,1), (1,-1), (-1,1), (-1,-1)} \rho(\lambda)d\lambda \\
&\int \rho(\lambda)d\lambda = 1 \quad \downarrow
\end{aligned}$$

$$|C(a, b) - C(a, b')| + |C(a', b) + C(a', b')| \leq 2$$

: Bell's inequality

This result contradicts QM prediction:

$$\begin{aligned}
C(a, b) &= \langle A(a)B(b) \rangle \quad \leftarrow \text{quantum average} \\
\theta_1 \uparrow \theta_2 \uparrow &= P(+, \theta_2; +, \theta_1) \times (+1) + P(-, \theta_2; -, \theta_1) \times (+1) \\
&\quad \quad \quad \nearrow \quad \quad \quad \nearrow \\
&\quad \quad \quad A(a)B(b) = (+1) \times (+1) \quad \quad \quad A(a)B(b) = (-1) \times (-1) \\
&+ P(+, \theta_2; -, \theta_1) \times (-1) + P(-, \theta_2; +, \theta_1) \times (-1) \\
&\quad \quad \quad \nearrow \quad \quad \quad \nearrow \\
&\quad \quad \quad A(a)B(b) = (-1) \times (+1) \quad \quad \quad A(a)B(b) = (+1) \times (-1)
\end{aligned}$$

$$\begin{aligned}
&= \alpha_1 \alpha_2 [\sin^2(\theta_1 - \theta_2) - \cos^2(\theta_1 - \theta_2)] \\
&= -\alpha_1 \alpha_2 \cos[2(\theta_1 - \theta_2)]
\end{aligned}$$

If $\alpha_1 = \alpha_2 = 1$ and $\theta_1 = 0, \theta_2 = \frac{3}{8}\pi, \theta'_1 = -\frac{\pi}{4}, \theta'_2 = \frac{\pi}{8}$, we obtain


$$\begin{aligned}
&|C(\theta_1, \theta_2) - C(\theta_1, \theta'_2)| + |C(\theta'_1, \theta_2) + C(\theta'_1, \theta'_2)| \\
&= \left| -\cos\left(-\frac{3}{4}\pi\right) + \cos\left(-\frac{\pi}{4}\right) \right| + \left| -\cos\left(-\frac{5}{4}\pi\right) - \cos\left(-\frac{3}{4}\pi\right) \right| \\
&= 2\sqrt{2} \quad \implies \text{violates Bell's inequality!}
\end{aligned}$$

This is a maximum violation point for Bell's inequality.

However, if $\alpha_1 \alpha_2 \leq \frac{1}{\sqrt{2}} \approx 0.71$, it is impossible to show the violation of Bell's inequality.

1.6.3 Clauser – Horne – Shimony - Holt (CHSH) inequality no enhancement assumption:

$$P(\theta, \lambda) \leq P(-, \lambda)$$



probability for a photon
to reach the detector
with a polarizer angle θ
probability for a photon
to reach the detector
without a polarizer

$$P_{12}(\theta_1, \theta_2) = \alpha_1 \alpha_2 \int P_1(\theta_1, \lambda) P_2(\theta_2, \lambda) \rho(\lambda) d\lambda$$

$$P_{12}(\theta_1, -) = \alpha_1 \alpha_2 \int P_1(\theta_1, \lambda) P_2(-, \lambda) \rho(\lambda) d\lambda$$

$$P_{12}(-, \theta_2) = \alpha_1 \alpha_2 \int P_1(-, \lambda) P_2(\theta_2, \lambda) \rho(\lambda) d\lambda$$

$$\left(\begin{array}{l} A(a, \lambda) \rightarrow P_1(\theta_1, \lambda) \\ B(b, \lambda) \rightarrow P_2(\theta_2, \lambda) \end{array} \right)$$

normalized probability:

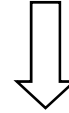
$$x = P_1(\theta_1, \lambda) / P_1(-, \lambda)$$

$$x' = P_1(\theta'_1, \lambda) / P_1(-, \lambda)$$

$$y = P_2(\theta_2, \lambda) / P_2(-, \lambda)$$

$$y' = P_2(\theta'_2, \lambda) / P_2(-, \lambda)$$

$$\implies 0 \leq x, x', y, y' \leq 1$$



$$\boxed{-1 \leq xy - xy' + x'y + x'y' - x' - y \leq 0}$$

$$\int \underbrace{\alpha_1 \alpha_2 P_1(-, \lambda) P_2(-, \lambda) \rho(\lambda)}_{\text{positive}} \underbrace{[xy - xy' + x'y + x'y' - x' - y]}_{\text{negative}} d\lambda$$

positive

negative



$$S = P_{12}(\theta_1, \theta_2) - P_{12}(\theta_1, \theta'_2) + P_{12}(\theta'_1, \theta_2) + P_{12}(\theta'_1, \theta'_2) - P_{12}(\theta'_1, -) - P_{12}(-, \theta_2) \leq 0$$



every term is proportional to $\alpha_1 \alpha_2$

$$\frac{S}{\alpha_1 \alpha_2} \leq 0 \quad \text{:CHSH inequality}$$

Since CHSH inequality is independent of $\alpha_1 \alpha_2$, a high-efficiency photo-detector is not necessary to demonstrate the violation of the inequality.

cf. QM prediction

$$\theta_1 = 0, \theta'_1 = -\frac{\pi}{4}, \theta_2 = \frac{3\pi}{8}, \theta'_2 = \frac{\pi}{8}$$

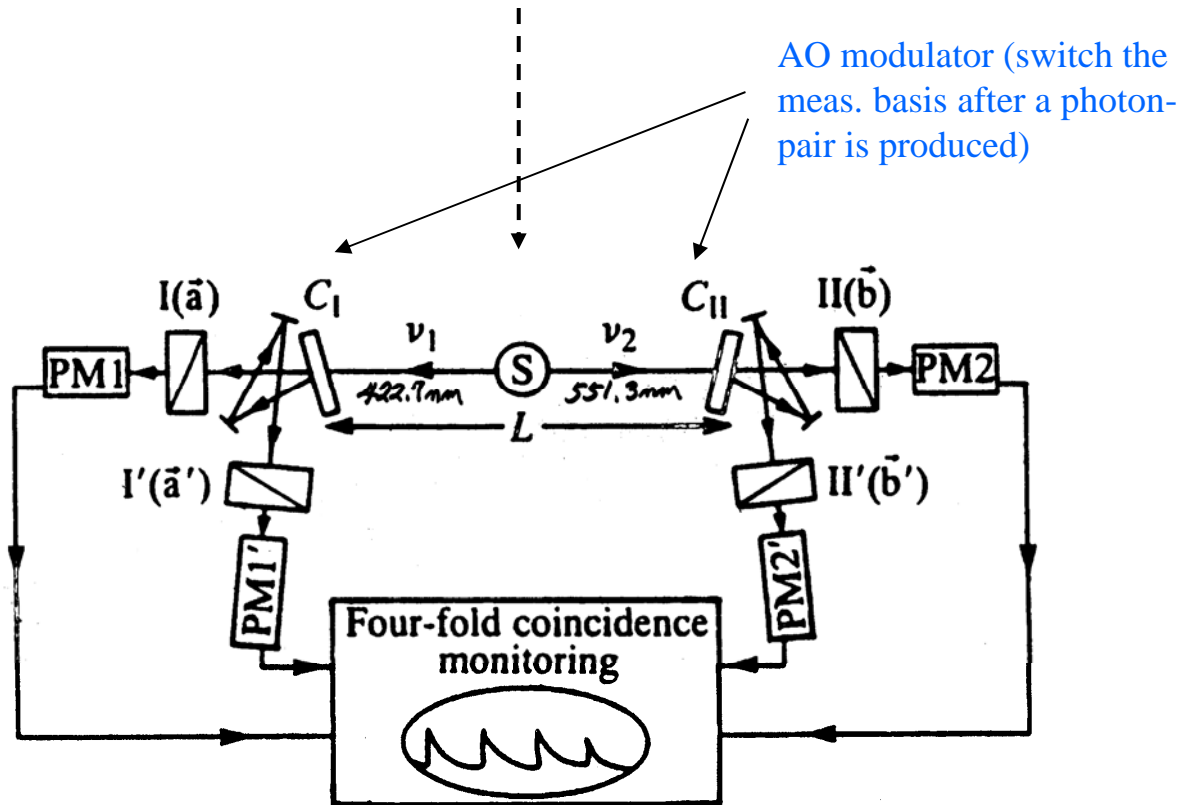


$$\frac{S}{\alpha_1 \alpha_2} = \frac{\sqrt{2} - 1}{2} > 0 \quad \text{(maximum violation of CHSH inequality)}$$

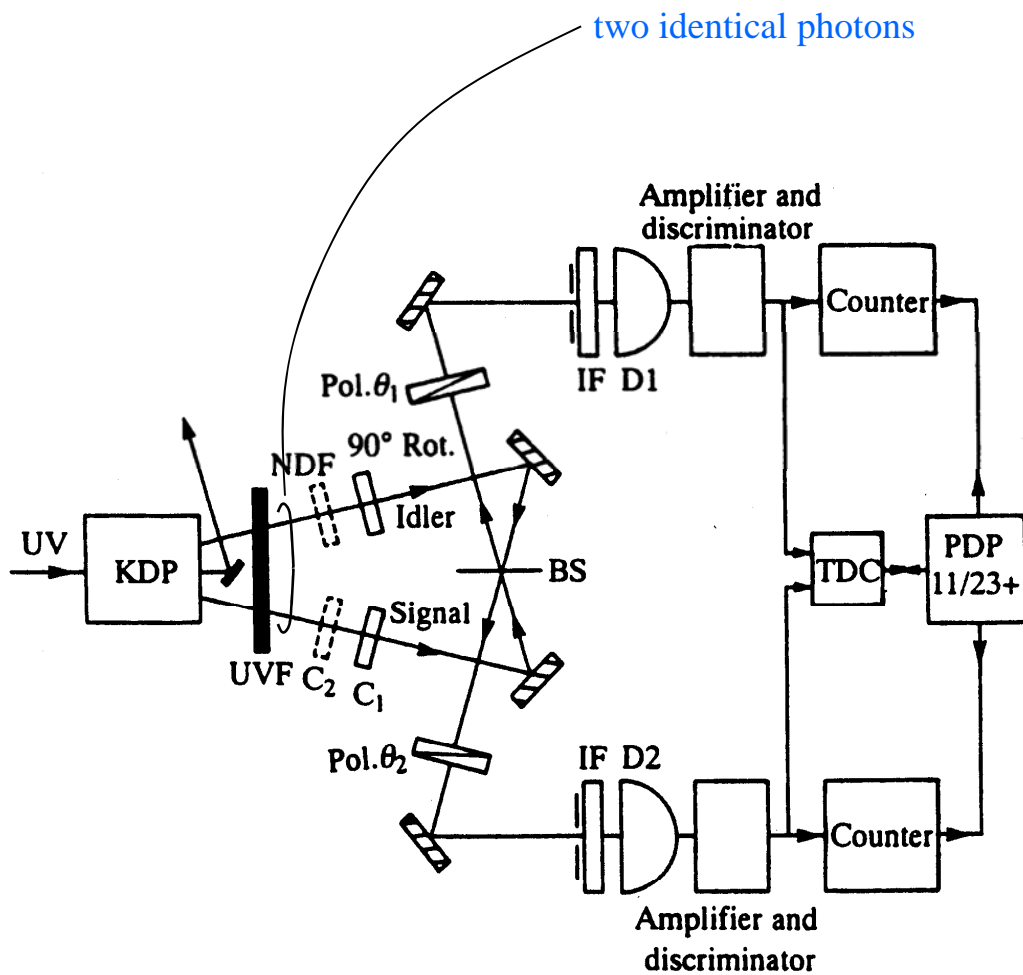
Thus, QM predicts the violation of the CHSH inequality.

Ca atomic SPS cascade emission

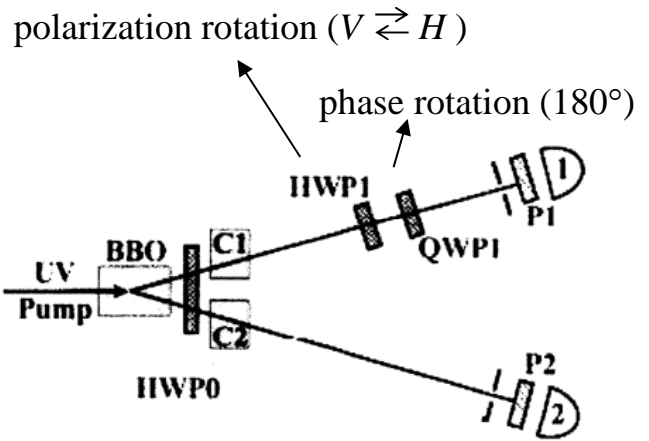
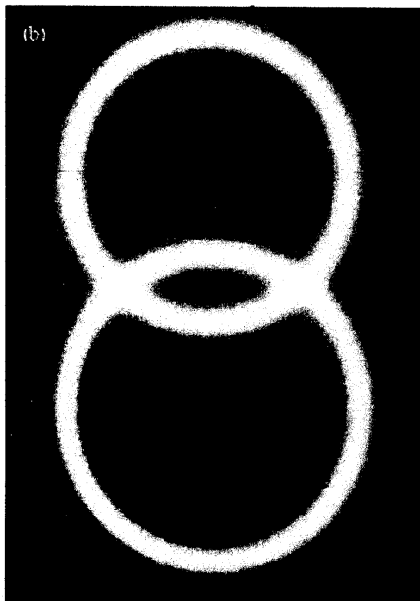
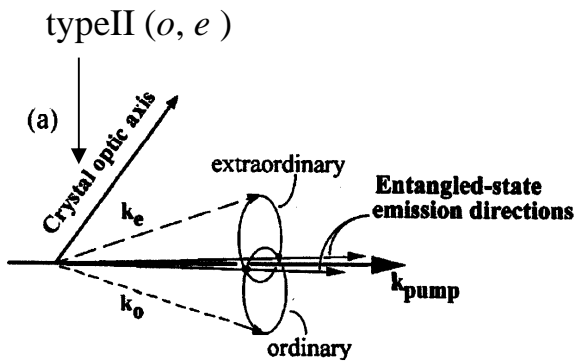
$$(J = 0) \rightarrow (J = 1) \rightarrow (J = 0)$$



A. Aspect, J. Dalibard and G. Roger, *Phys. Rev. Lett.* **49**, 1804 (1982)



Z. Y. Ou and L. Mandel, *Phys. Rev. Lett.* **61**, 50 (1988)

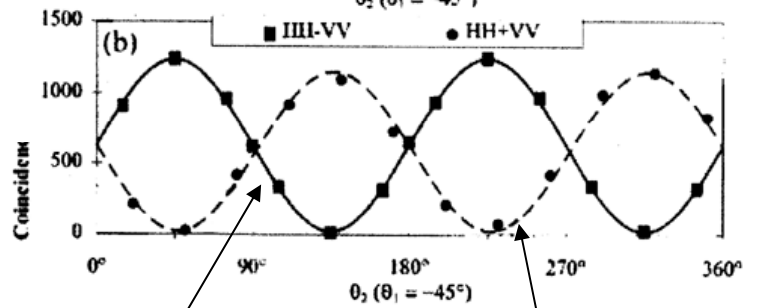
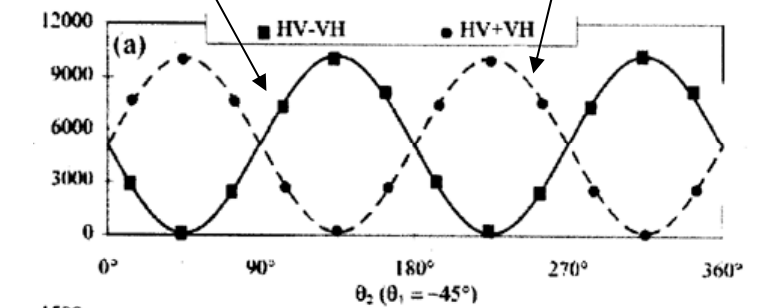


$$\frac{1}{\sqrt{2}} [|H_1V_2\rangle - |V_1H_2\rangle]$$

$$\sin^2(\theta_1 - \theta_2)$$

$$\frac{1}{\sqrt{2}} [|H_1V_2\rangle + |V_1H_2\rangle]$$

$$\sin^2(\theta_1 + \theta_2)$$



$$\frac{1}{\sqrt{2}} [|H_1H_2\rangle - |V_1V_2\rangle]$$

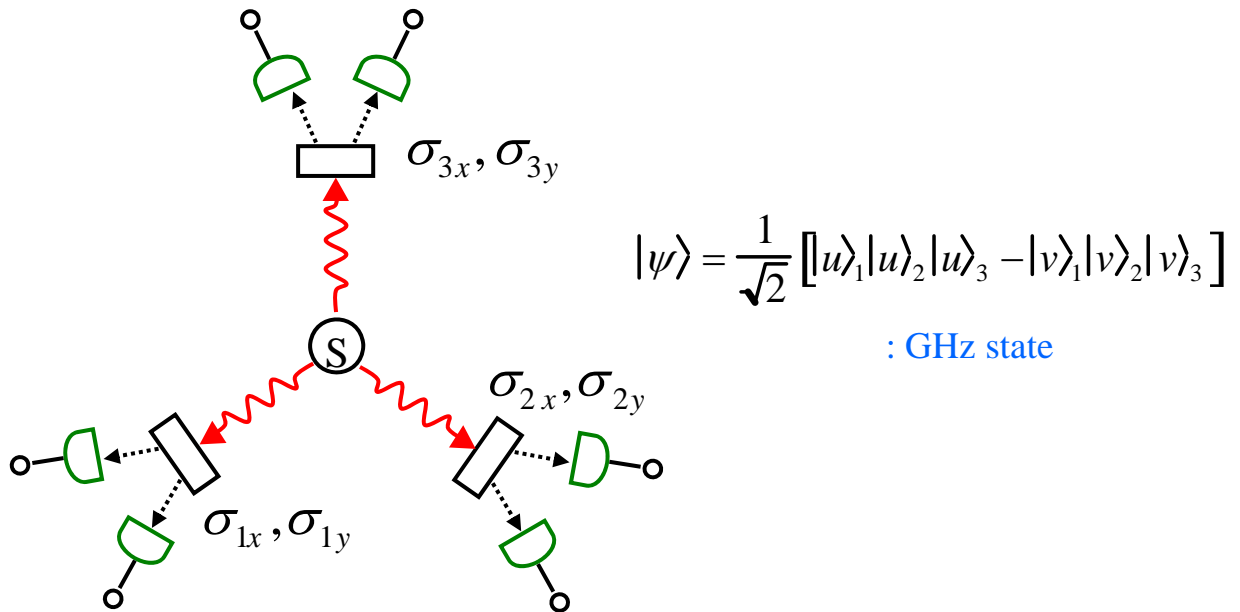
$$\cos^2(\theta_1 + \theta_2)$$

$$\frac{1}{\sqrt{2}} [|H_1H_2\rangle + |V_1V_2\rangle]$$

$$\cos^2(\theta_1 - \theta_2)$$

P. G. Kwiat, K. Mattle, H. Weinfurter and A. Zeilinger,
Phys. Rev. Lett. **75**, 4337 (1995)

1.6.4 Three Particle Entangled State (spin-1/2 particles)



Projective properties of Pauli spin operators:

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\sigma_y = \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix}$$

$$\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\sigma_x |u\rangle = |v\rangle$$

$$\sigma_y |u\rangle = i|v\rangle$$

$$\sigma_z |u\rangle = |u\rangle$$

$$\sigma_x |v\rangle = |u\rangle$$

$$\sigma_y |v\rangle = -i|u\rangle$$

$$\sigma_z |v\rangle = -|v\rangle$$

$|u\rangle \pm |v\rangle$ are the eigenstates
of σ_x with eigenvalues of ± 1

$|u\rangle \pm i|v\rangle$ are the eigenstates
of σ_y with eigenvalues of ± 1

Eigenvalue properties of $|\psi\rangle$:

$$\underbrace{\sigma_{1x} \sigma_{2y} \sigma_{3y}}_{\text{projective meas. of } 1x, 2y, 3y} |\psi\rangle = \frac{1}{\sqrt{2}} [(i)^2 |v\rangle_1 |v\rangle_2 |v\rangle_3 - (-i)^2 |u\rangle_1 |u\rangle_2 |u\rangle_3] = |\psi\rangle$$

\Rightarrow (xyy) measurement

$$\sigma_{1y} \sigma_{2x} \sigma_{3y} |\psi\rangle = |\psi\rangle$$

$$\sigma_{1y} \sigma_{2y} \sigma_{3x} |\psi\rangle = |\psi\rangle$$

$$\sigma_{1x} \sigma_{2x} \sigma_{3x} |\psi\rangle = \frac{1}{\sqrt{2}} [|v\rangle_1 |v\rangle_2 |v\rangle_3 - |u\rangle_1 |u\rangle_2 |u\rangle_3] = -|\psi\rangle$$

Measurement results:

$m_{1x}m_{2x}m_{3x} = -1$ ← If the spins of the particles #2 and #3 are measured along x-axis, the measurement result of the spin of the particle #1 along x-axis is predicted with certainty.

4 different measurements {

- 1 ← (1, 1)
- 1 ← (1, -1)
- 1 ← (-1, 1)
- 1 ← (-1, -1)

$m_{1x}m_{2y}m_{3y} = 1$ ← If the spins of the particles #2 and #3 are measured along y-axis, the measurement result of the spin of the particle #1 along x-axis is predicted with certainty.

- 1 ← (1, 1)
- 1 ← (1, -1)
- 1 ← (-1, 1)
- 1 ← (-1, -1)

The above two m_{1x} 's are identical according to Einstein's local realism

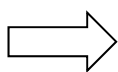
← An element of physical reality

$m_{1y}m_{2x}m_{3y} = 1$

$m_{1y}m_{2y}m_{3x} = 1$



$(m_{1x}m_{1y}m_{2x}m_{2y}m_{3x}m_{3y})^2 = -1$: contradiction



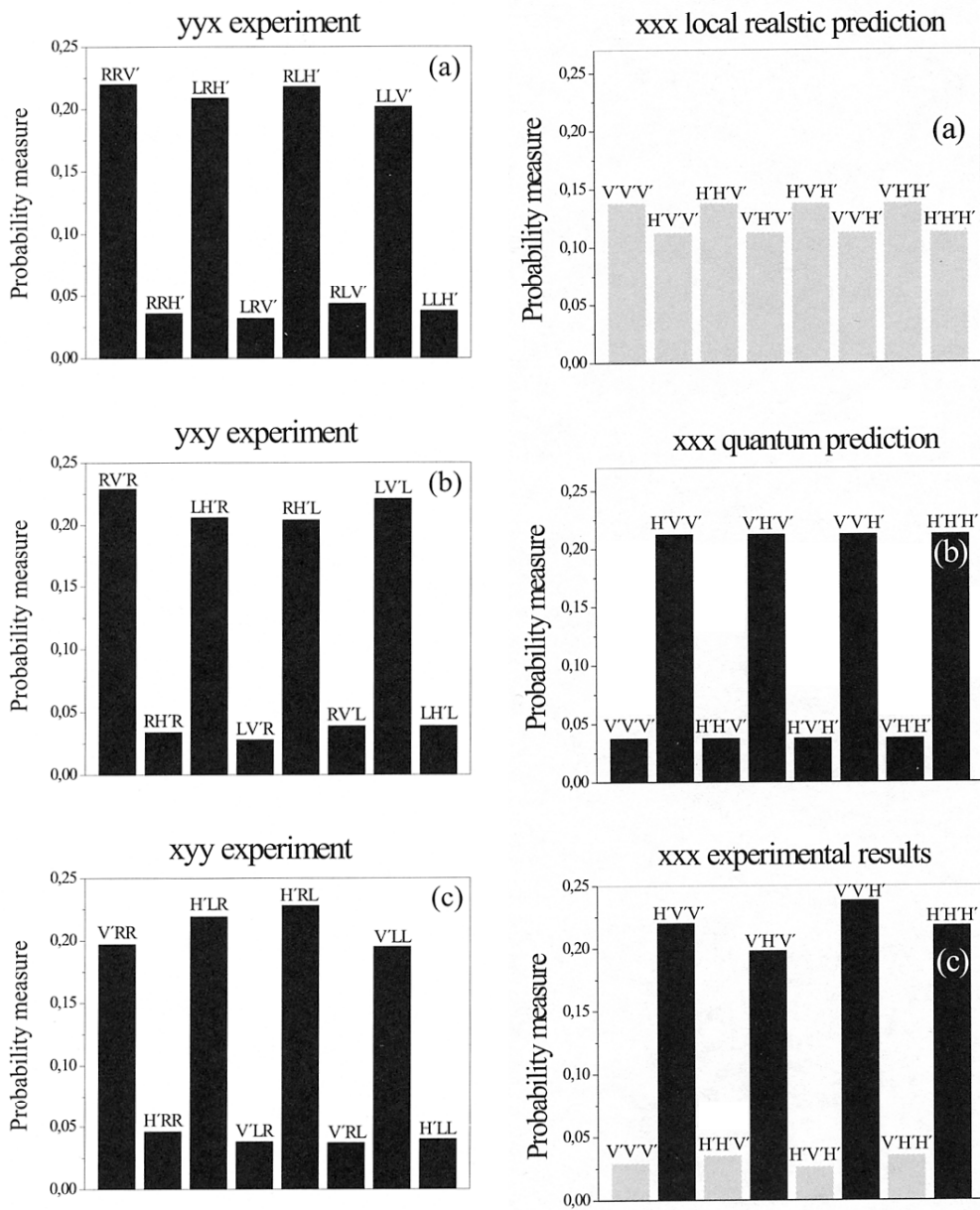
Violation of the local hidden variable theory without inequality

Experimental Test of Local Realism vs. Quantum Mechanics

GHZ state $|\psi_{12}\rangle = \frac{1}{\sqrt{2}} [|H\rangle_1 |H\rangle_2 |H\rangle_3 + |V\rangle_1 |V\rangle_2 |V\rangle_3]$

$|R\rangle = \frac{1}{\sqrt{2}} [|H\rangle + i|V\rangle]$ $|L\rangle = \frac{1}{\sqrt{2}} [|H\rangle - i|V\rangle]$

$|H'\rangle = \frac{1}{\sqrt{2}} [|H\rangle + |V\rangle]$ $|V'\rangle = \frac{1}{\sqrt{2}} [|H\rangle - |V\rangle]$



Further reading for Chapter 1

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